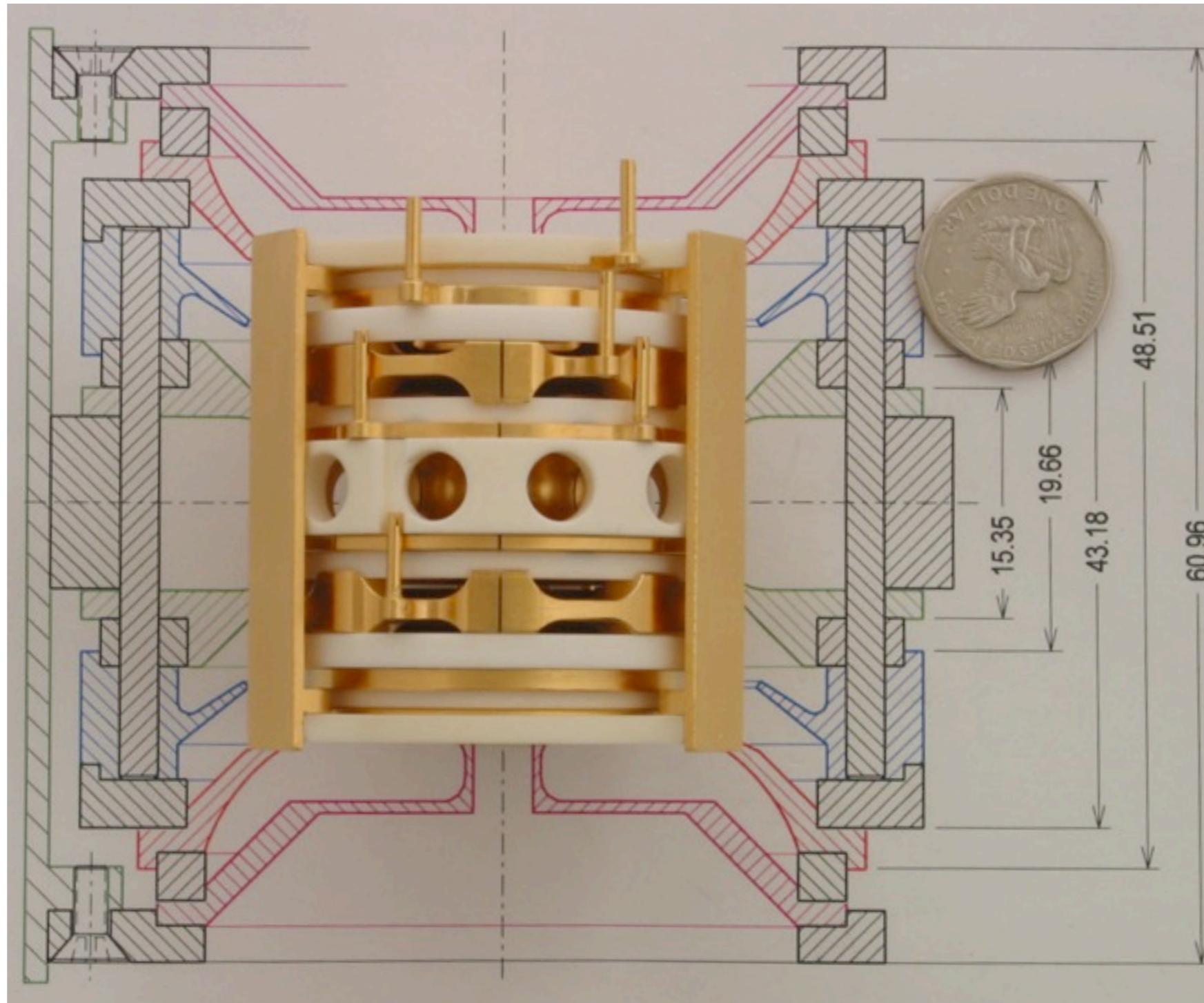
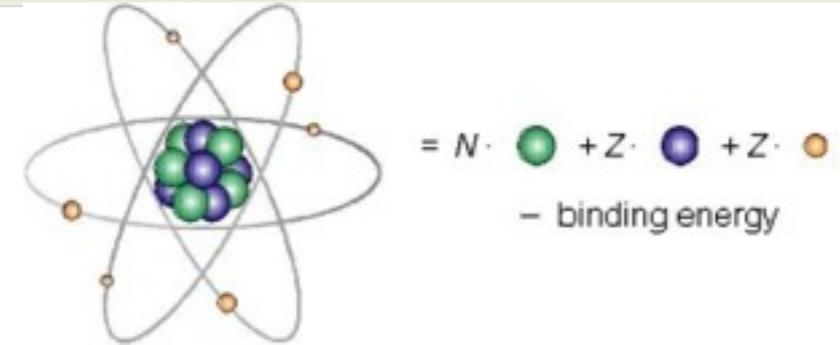


Plans for further technical development and improvements in trap-based mass measurements



Outline

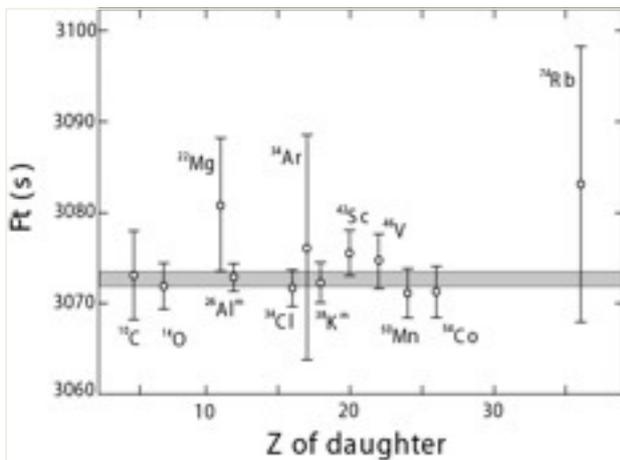
- Range of relative mass precisions required for mass measurements of rare isotopes
- Principles of Penning trap operation
 - confining fields
 - multipolar RF fields to drive ion motion
 - time-of-flight resonant detection technique
 - mass resolution
- State of the art and future technical developments
 - Short and long term implementation
 - NSCL's possible contributions



Masses of rare isotopes

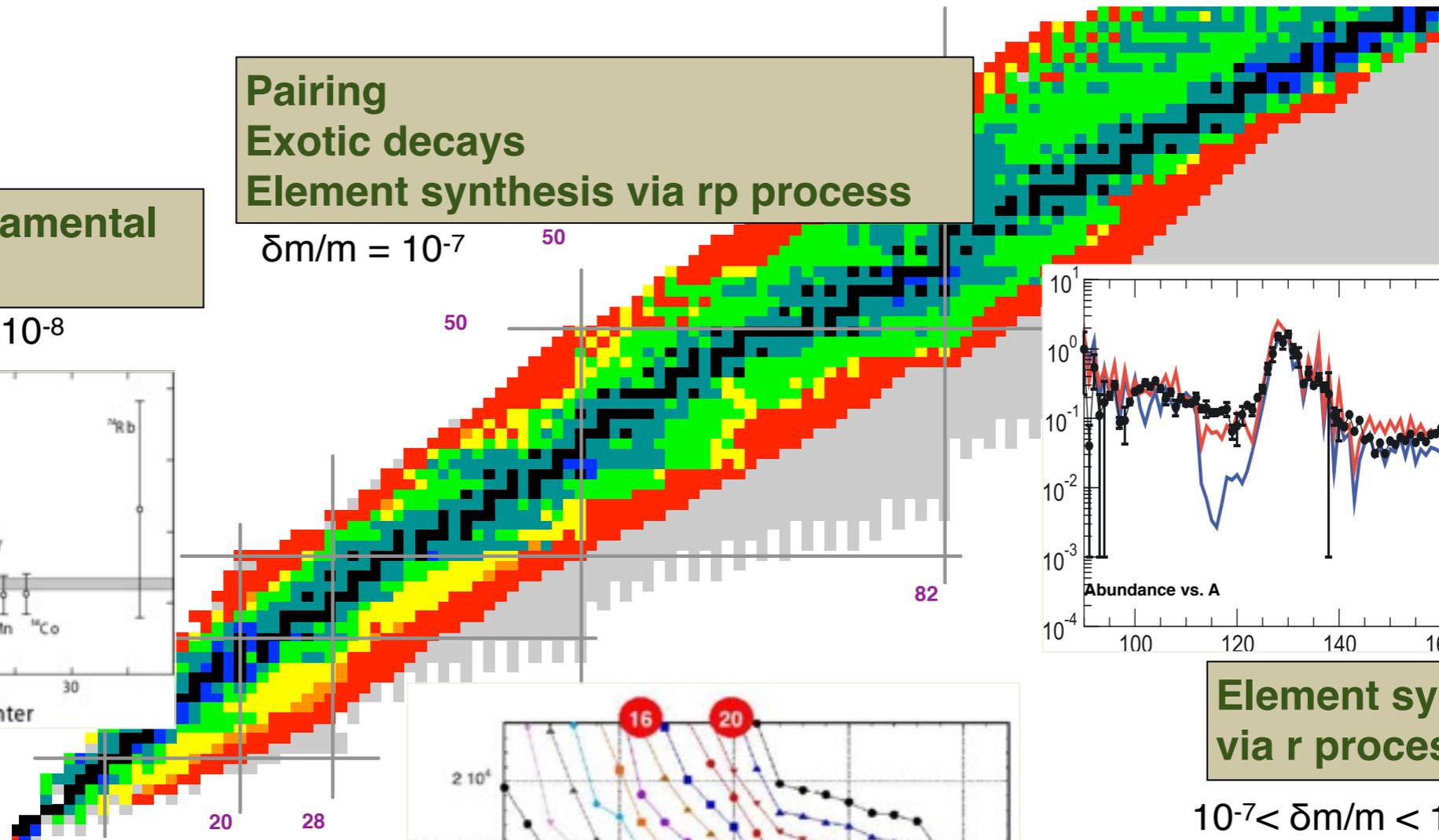
Tests of Fundamental Interactions

$$\delta m/m < 10^{-8}$$



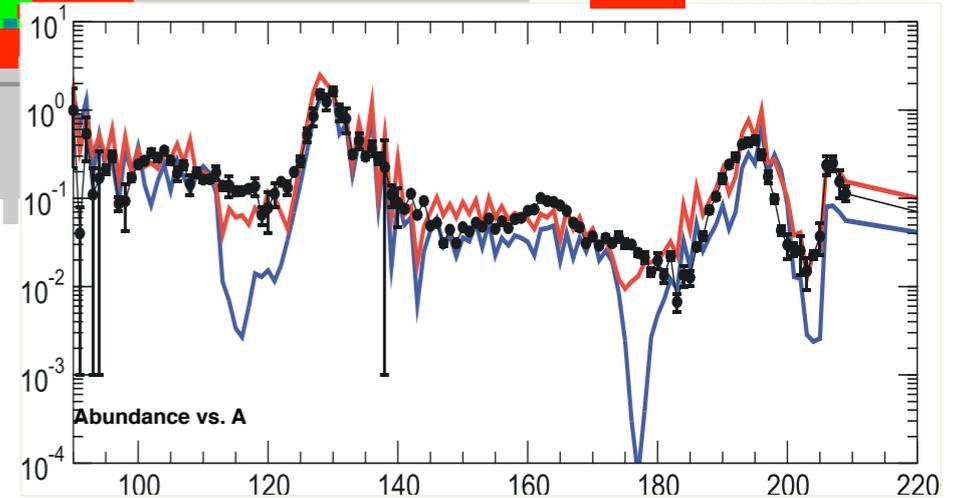
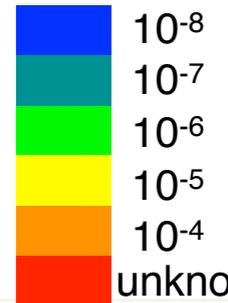
Pairing
Exotic decays
Element synthesis via rp process

$$\delta m/m = 10^{-7}$$



Mass Uncertainties

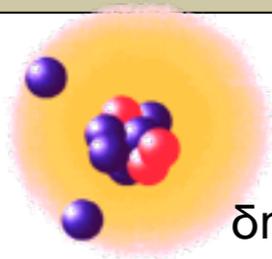
$\delta m/m$



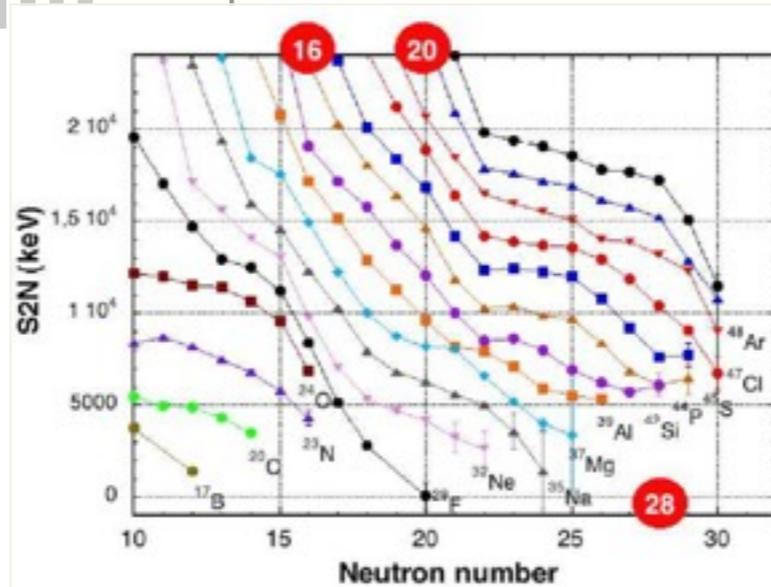
Element synthesis via r process

$$10^{-7} < \delta m/m < 10^{-6}$$

Halos and skins



$$\delta m/m = 10^{-7}$$



Evolution of nuclear shell structure

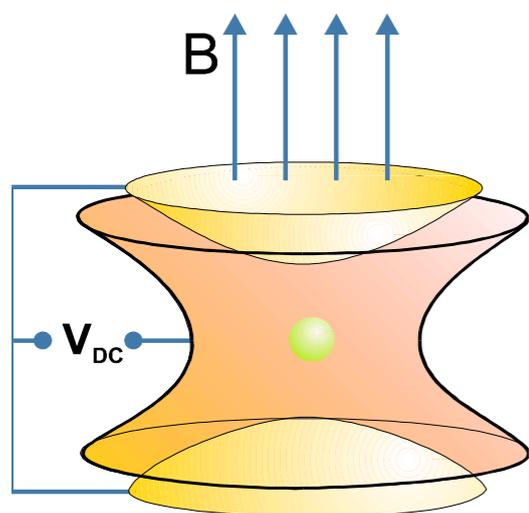
$$10^{-6} < \delta m/m < 10^{-5}$$



Penning trap basics

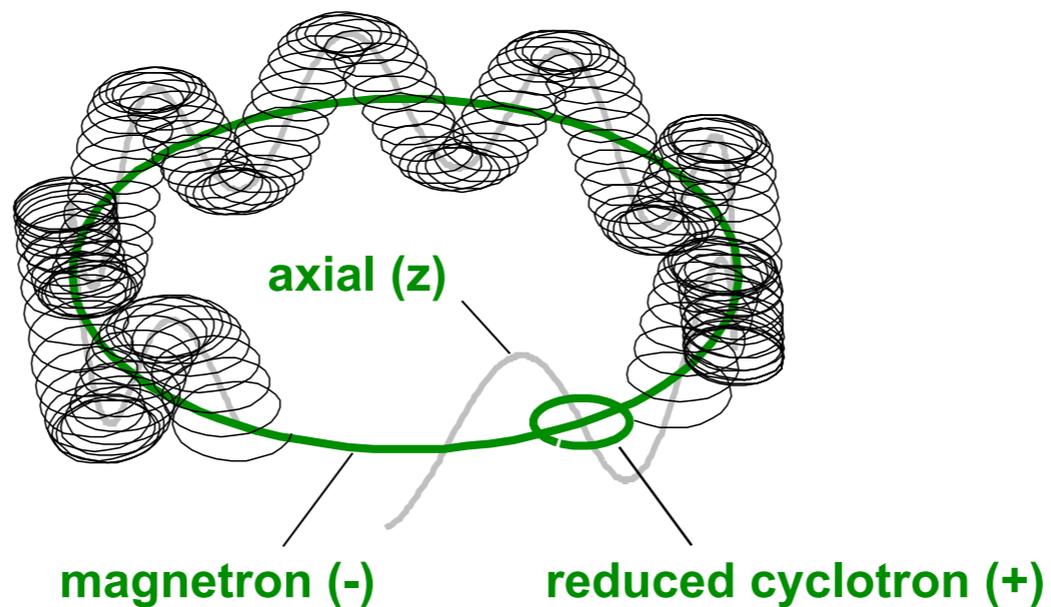


homogenous magnetic field



electrode structure

trapped ions execute three independent eigenmotions^{1,2}



Important relation:

$$\nu_c = \frac{q}{2\pi m} \cdot B = \nu_+ + \nu_-$$

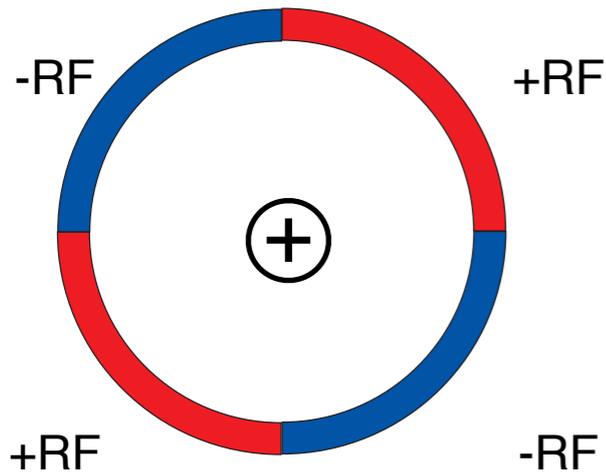
eigenmotions can be driven with the application of RF fields

- mass measurements
- isobaric/isomeric purification
- cooling/centering of beam

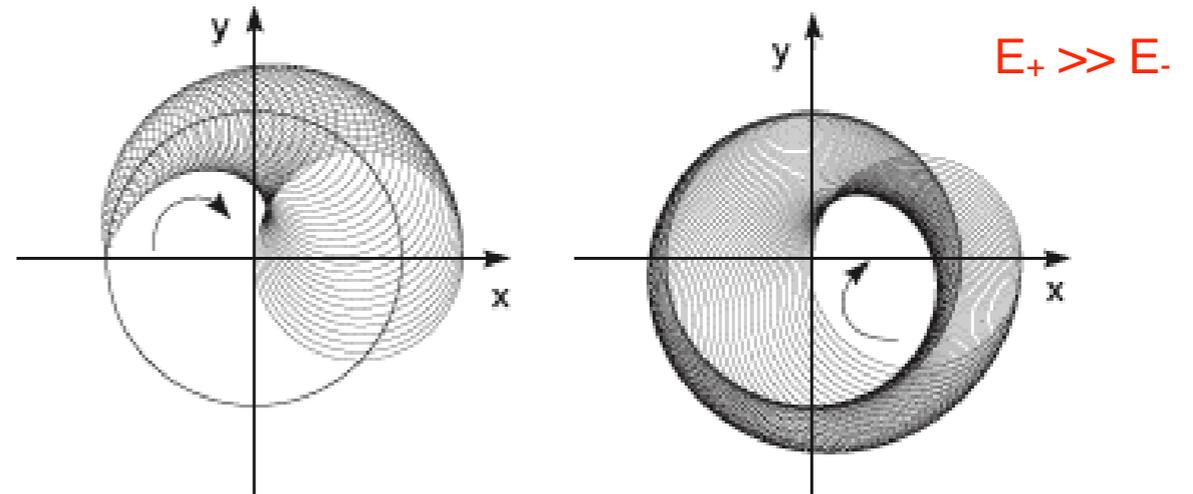
1. L. S. Brown and G. Gabrielse, *Reviews of Modern Physics* **58**, 233-311 (1986).

2. M. König *et al.*, *Int. J. Mass Spectrom.* **142**, 95-116 (1995).

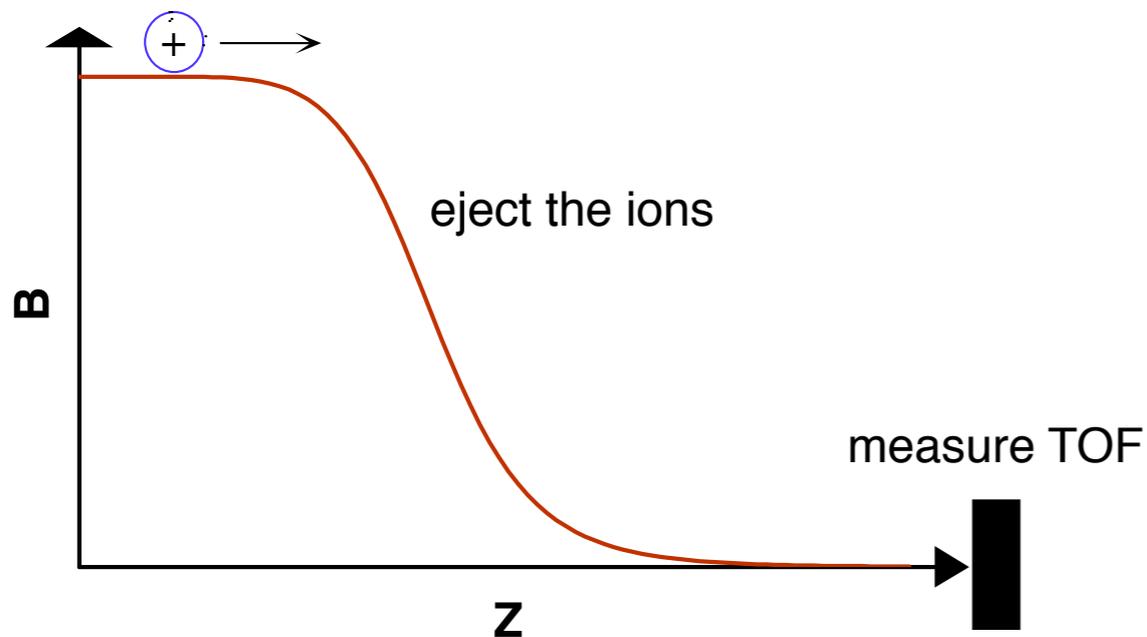
Quadrupolar excitation and resonance time-of-flight detection



trap ions and apply RF excitations



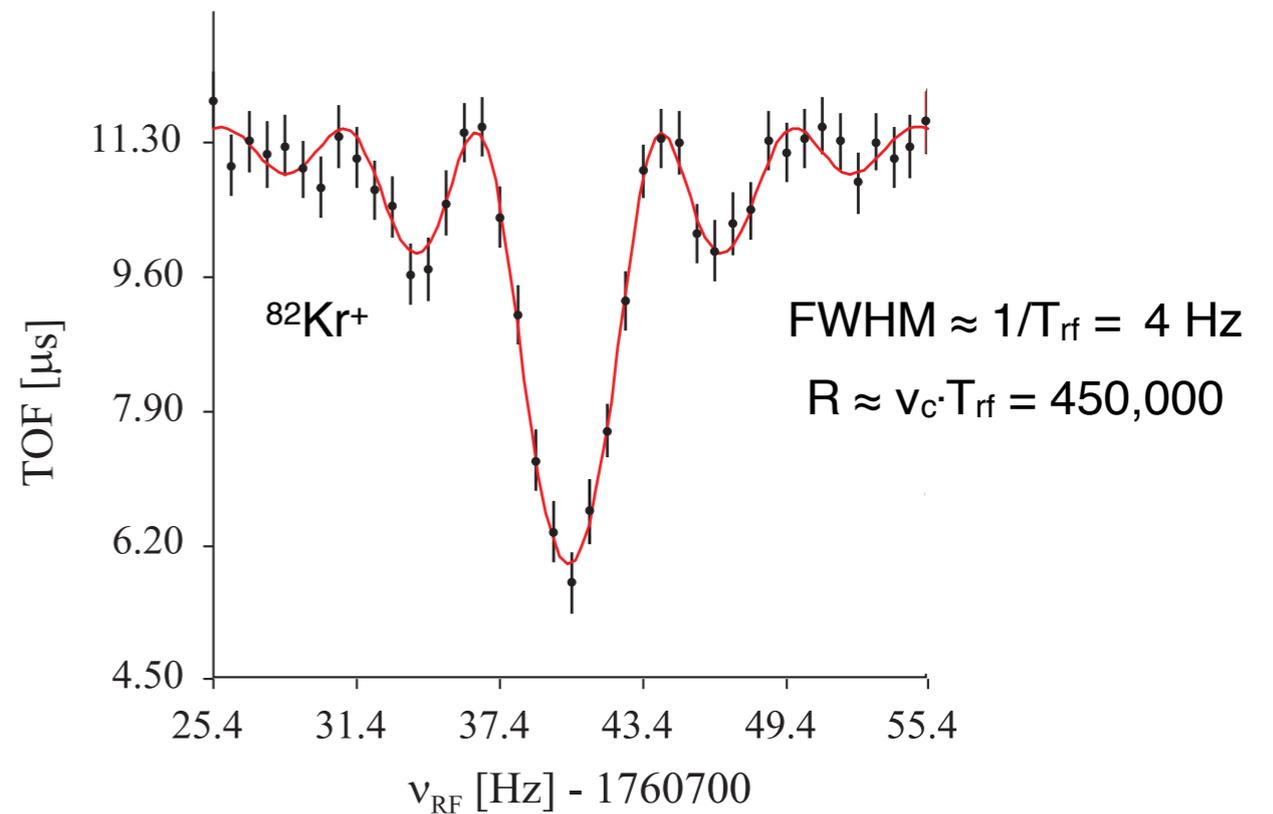
quadrupolar excitation: beating of motions when $\nu_{rf} \approx \nu_c$



eject the ions

measure TOF

$$F = -\mu \frac{\partial B}{\partial z} = -\frac{E_r}{B_0} \cdot \frac{\partial B}{\partial z}$$



What goes into relative mass precision?

$$\frac{\delta m}{m} = \frac{\gamma}{R \cdot \sqrt{N}}$$

γ - system-specific scaling factor

- initial ion distribution
- extraction conditions
- length of flight path
- contaminant ions

optimized injection
optimized ejection
efficient purification

R - resolving power of excitation

- function of T_{rf} and v_c

improved with:



better excitation schemes
charge breeding

N - number of detected ions

- detection efficiency
- yield, total experiment time
- measurement overhead

large duty cycle
fast purification

Fast, efficient purification

Isobaric purification of contaminant species

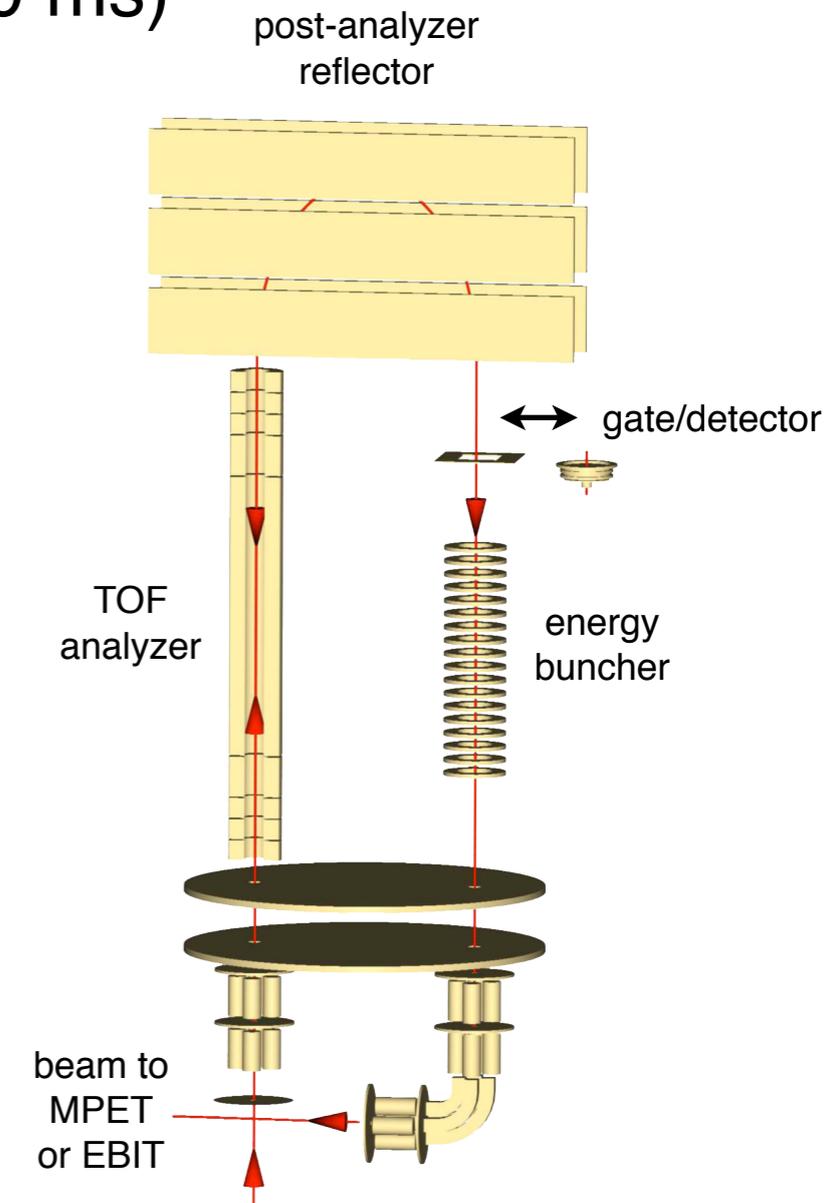
- sideband cooling in a gas-filled Penning trap¹ ($t \sim 100$ ms)
(broadband, no knowledge of contaminant required)
- dipolar excitation of contaminants ($t \sim 10$ ms)
(contaminant species need to be identified)

MR-TOF-MS Isobar Separator²

Principle: electrostatic mirror system drastically increases the ion flight path

Advantages:

1. extremely short measurement times (100 ns to 10 ms)
2. broad mass range
3. large ion capacity
4. high resolving power ($m/\Delta m \sim 100,000$)
5. compact setup



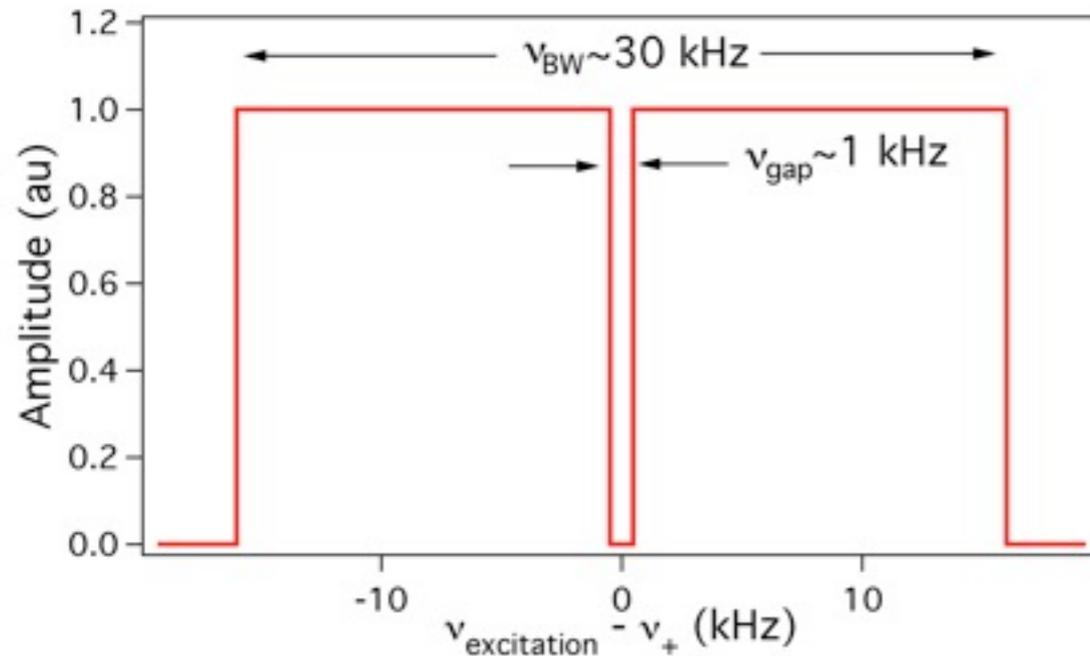
1. G. Bollen *et al.*, *J. Appl. Phys.* **68**, 4355-74 (1990).

2. W. R. Plaß *et al.*, *Nucl. Instrum. Methods Phys. Res., Sect. B* **266**, 4560-4 (2008).

Fast, efficient purification

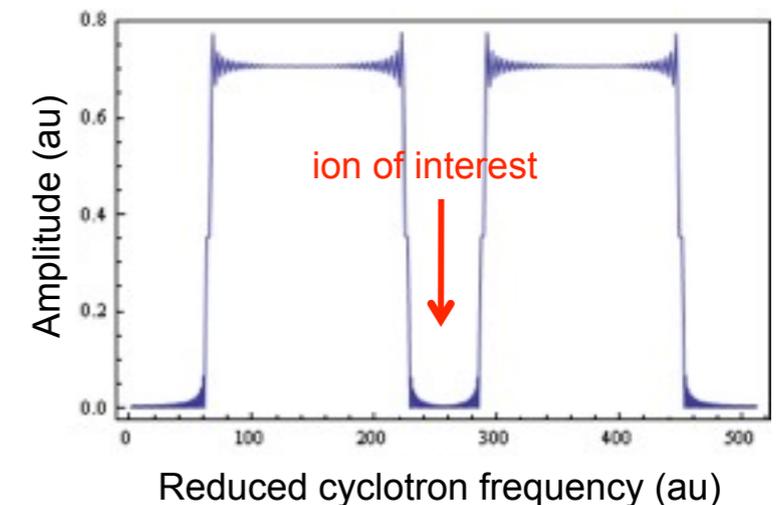
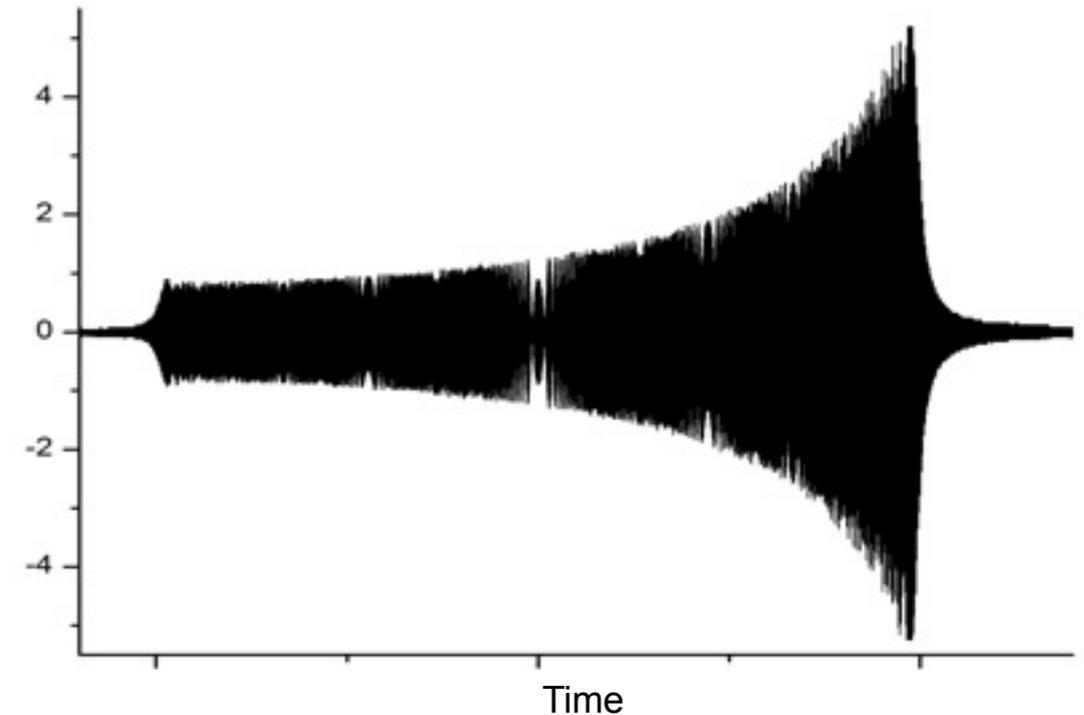
Implement SWIFT¹ Technique used in FT-ICR

Frequency domain signal



IFT
Amplitude (au)

Time domain signal



Requirements for LEBIT

- Isobaric Purification Molecular Isobars \Rightarrow Mass bandwidth ± 0.1 u
Nuclear Isobars $\Rightarrow \Delta m/m \sim 10^{-4}$

Challenge

- Minimize power leakage into gap

Requirements

- Programming, function generator, amplifier

SWIFT is a cheap, efficient and fast solution

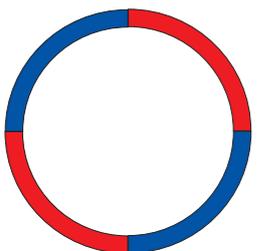
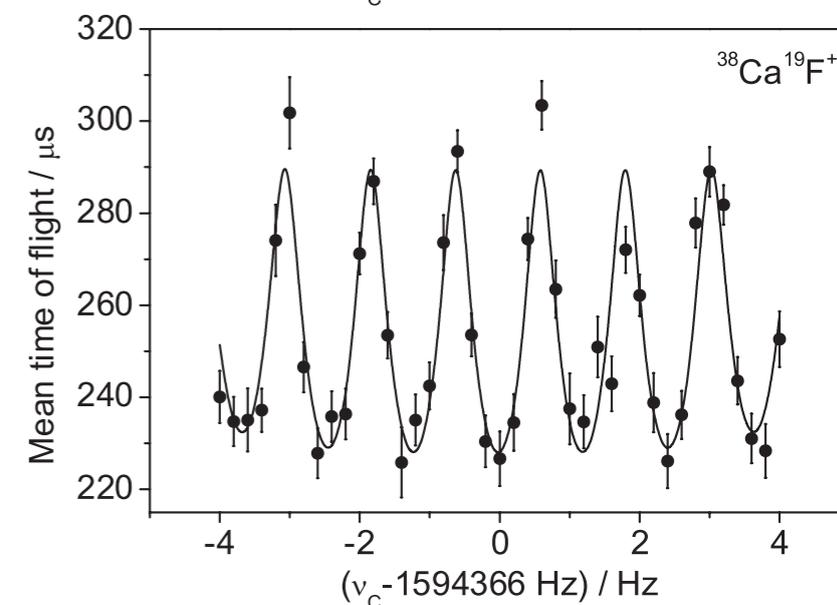
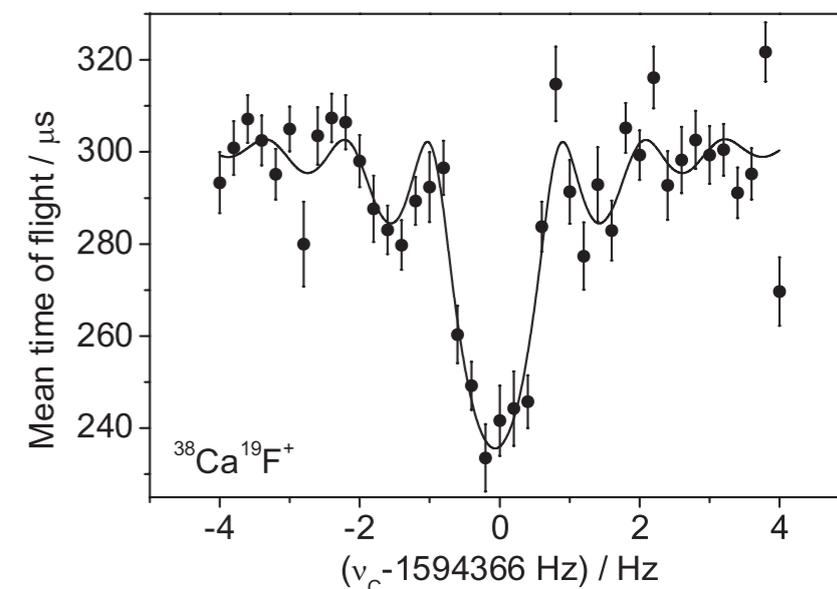
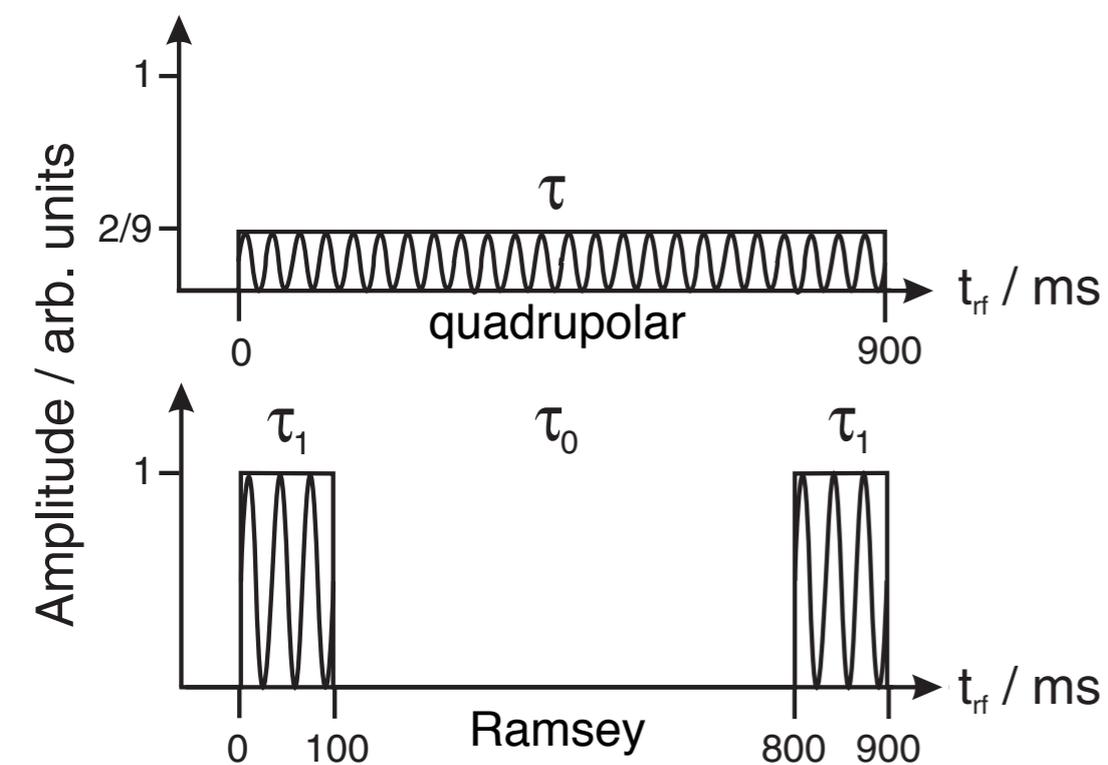
1. S. Guan and A. Marshall, *Int. J. Mass Spectrom.*, 5-37 (1996).



Excitation schemes

(using quadrupolar RF field)

Goal: highest resolving power for given T_{rf}



Quadrupolar: $R \sim T_{rf} \cdot \nu_c$

Ramsey^{1,2,3}: $R \sim 3 \cdot T_{rf} \cdot \nu_c$

Ramsey offers greater resolving power free of charge

- issues with isomer resolution, center freq. determination

1. N. F. Ramsey, *Reviews of Modern Physics* **62**, 541-52 (1990).
2. S. George *et al.*, *Phys. Rev. Lett.* **98**, 162501 (2007).
3. M. Kretschmar, *Int. J. Mass Spectrom.* **264**, 122-45 (2007).

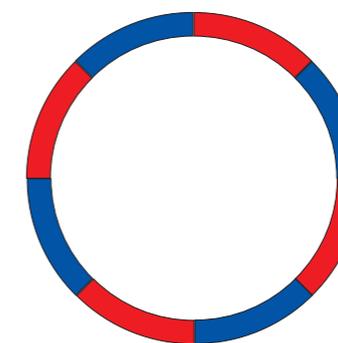


Excitation scheme

(using octupolar RF field)

requires eight-fold ring segmentation

$$v_{rf} \approx 2v_c$$



quadrupolar excitation:

- FWHM $\approx 1/T_{rf}$

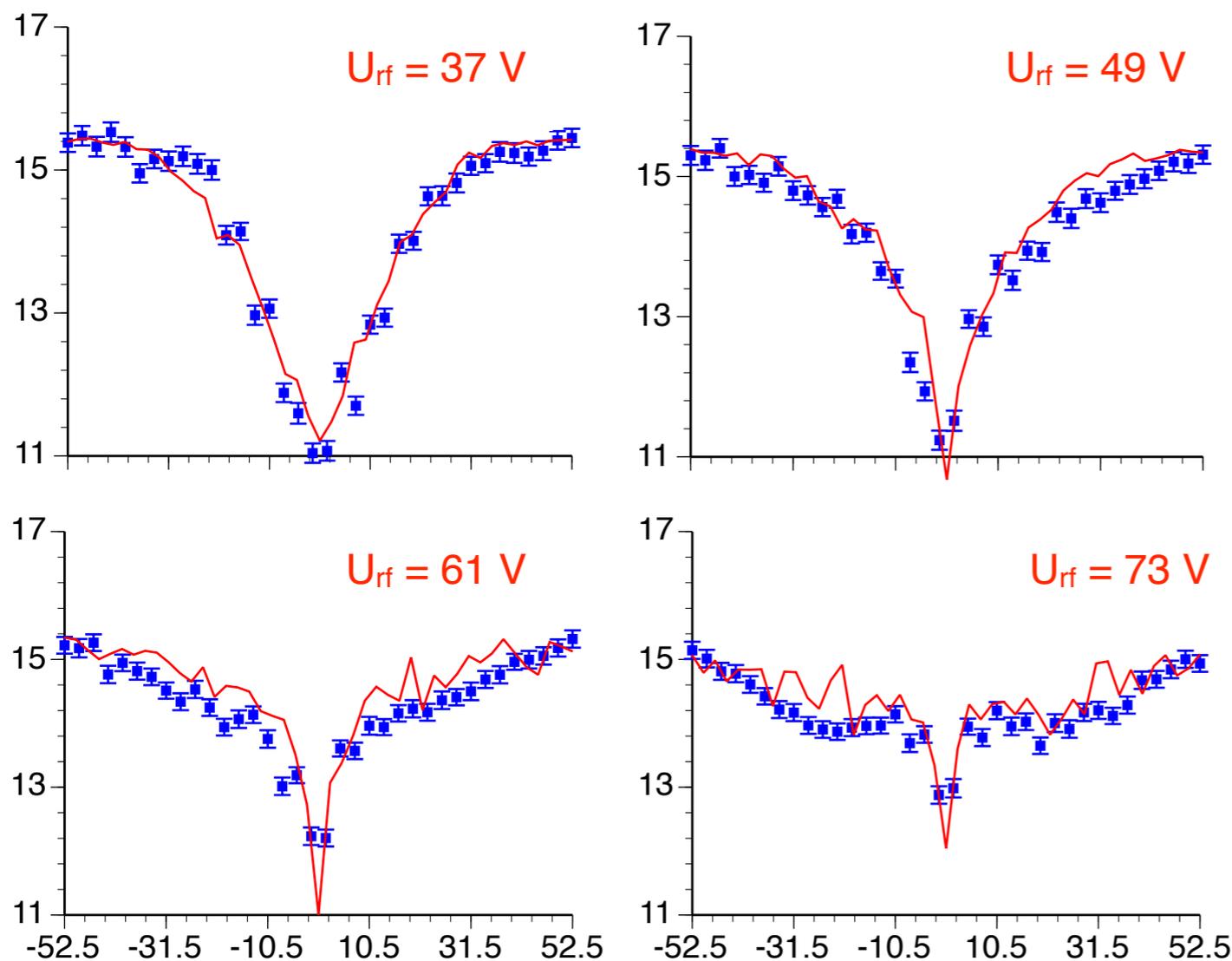
octupolar excitation^{1,2}:

- FWHM $\approx 0.2/T_{rf}$ @ $U_{rf} = 73$ V

$$\longrightarrow R_{oct} = 10 \cdot R_{quad}$$

Note:

ultimate gain in R is highly dependent on initial ion cloud distribution



$$v_{rf} = 12539572.5 \text{ [Hz]}$$

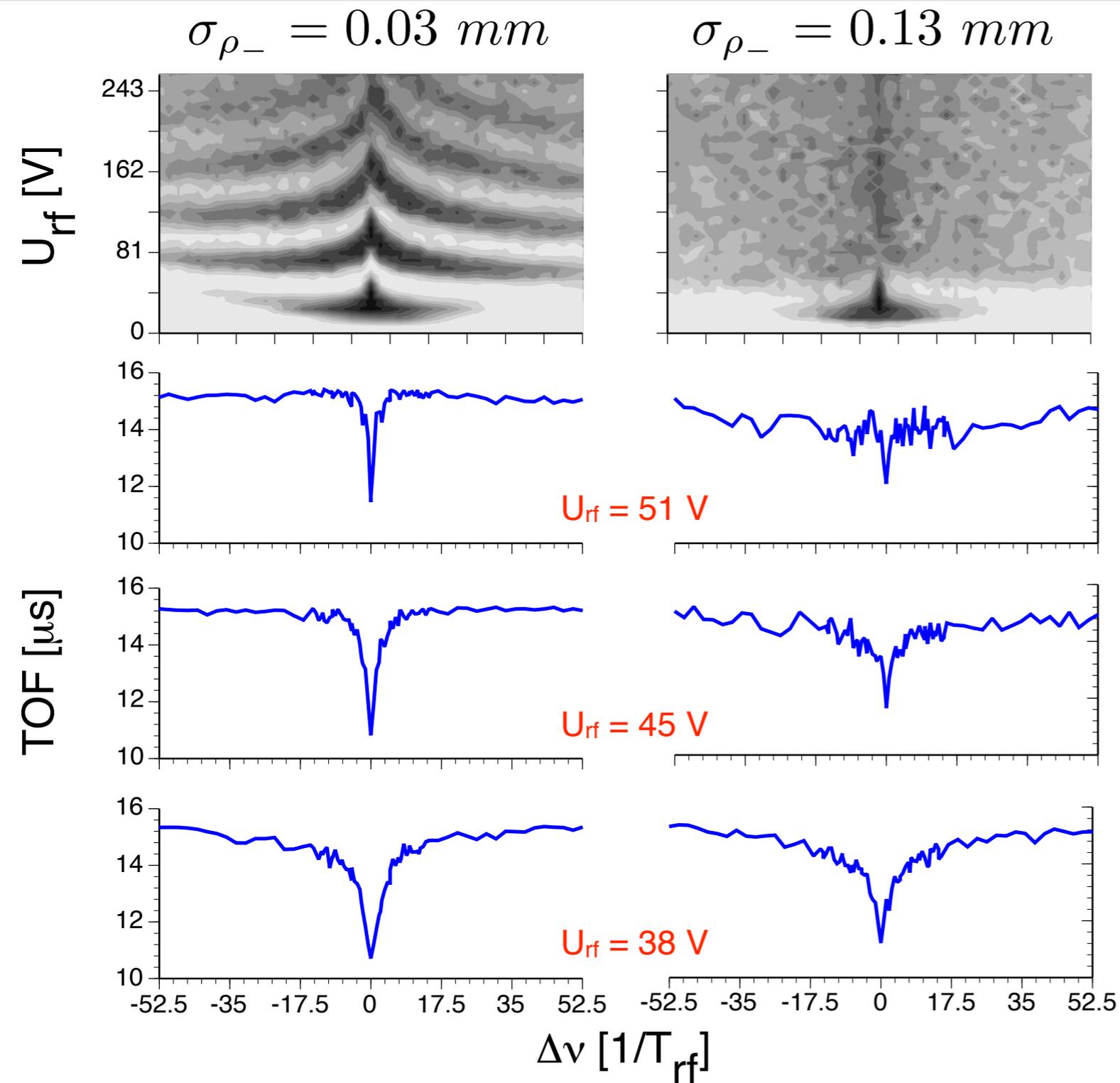
$$^{23}\text{Na}^+, T_{rf} = 50 \text{ ms}$$

1. R. Ringle *et al.*, *Int. J. Mass Spectrom.* **262**, 33-44 (2007).

2. S. Eliseev *et al.*, *Int. J. Mass Spectrom.* **262**, 45 - 50 (2007).

Excitation scheme

(using octupolar RF field)



Realistic multi-ion simulations:

- conversion frequency is dependent on amplitude
- smaller magnetron distribution yields higher resolving power
- increases in resolving powers of ≈ 20 within reach of current system
- preliminary results, further studies required
- no theoretical line shape

Experimental results:

- accuracy of $\sim 5 \times 10^{-9}$ experimentally verified
- $R_{\text{oct}} \sim 10 \cdot R_{\text{quad}}$

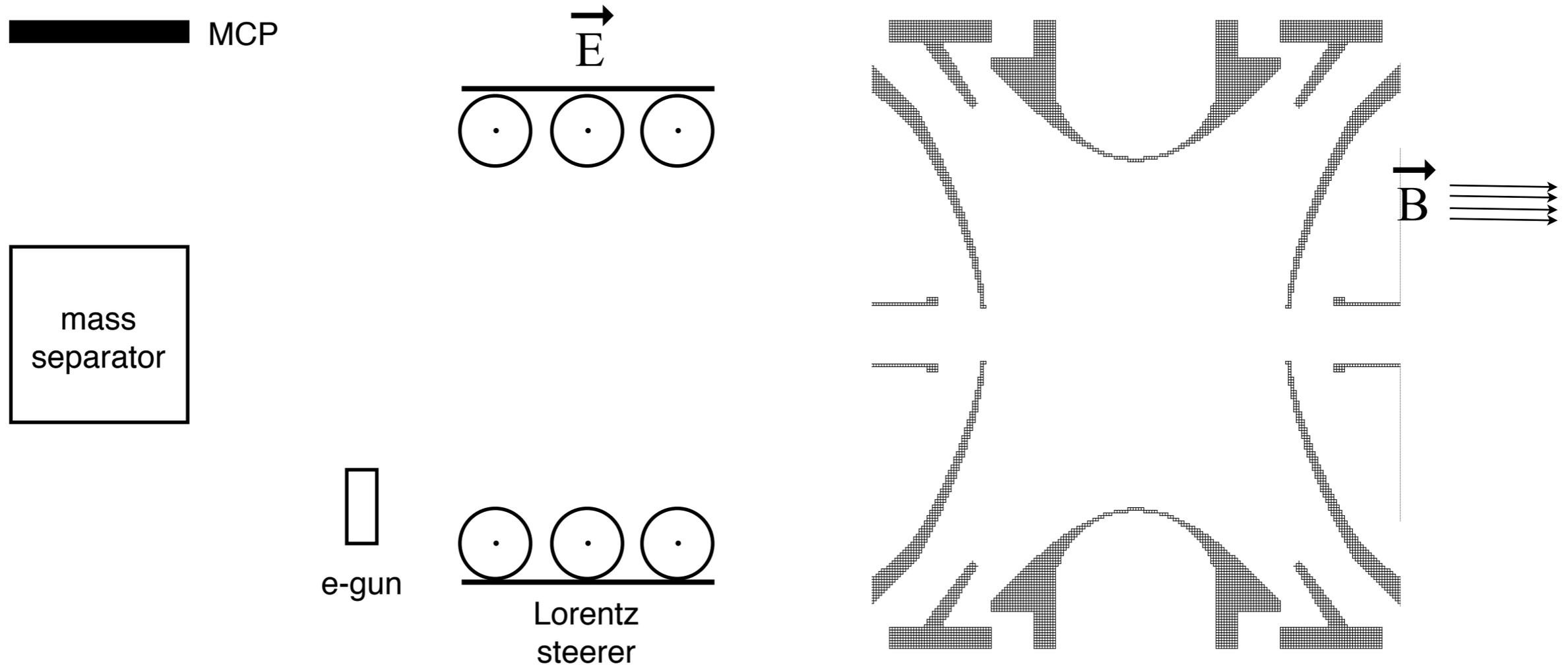
Charge breeding

Moderate charge breeding on a budget

- no EBIT and separate cooler trap required
- efficient use of all charge states produced
- less complicated

$$\nu_c = \frac{q}{2\pi m} \cdot B = \nu_+ + \nu_-$$

Increasing q increases ν_c
larger ν_c increases R for given T_{rf}



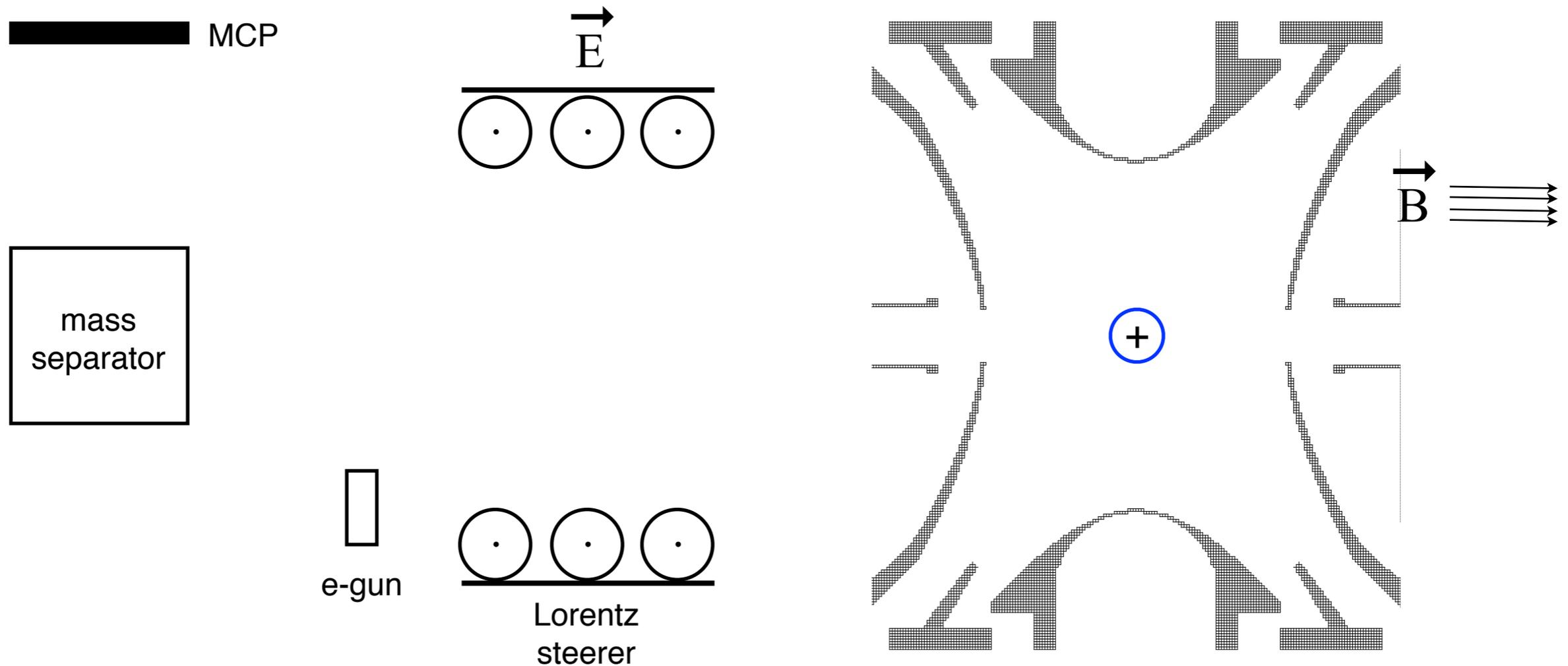
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Increasing q increases ν_c
 larger ν_c increases R for given T_{rf}



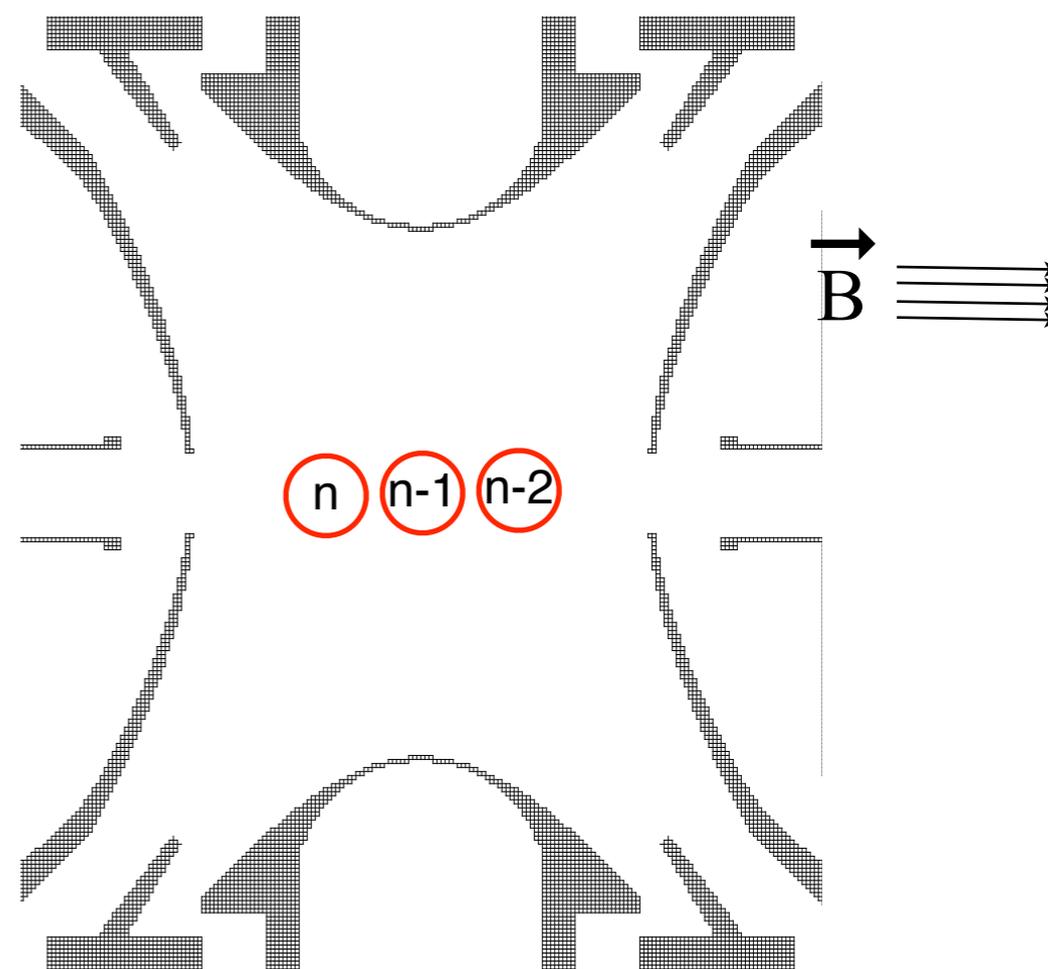
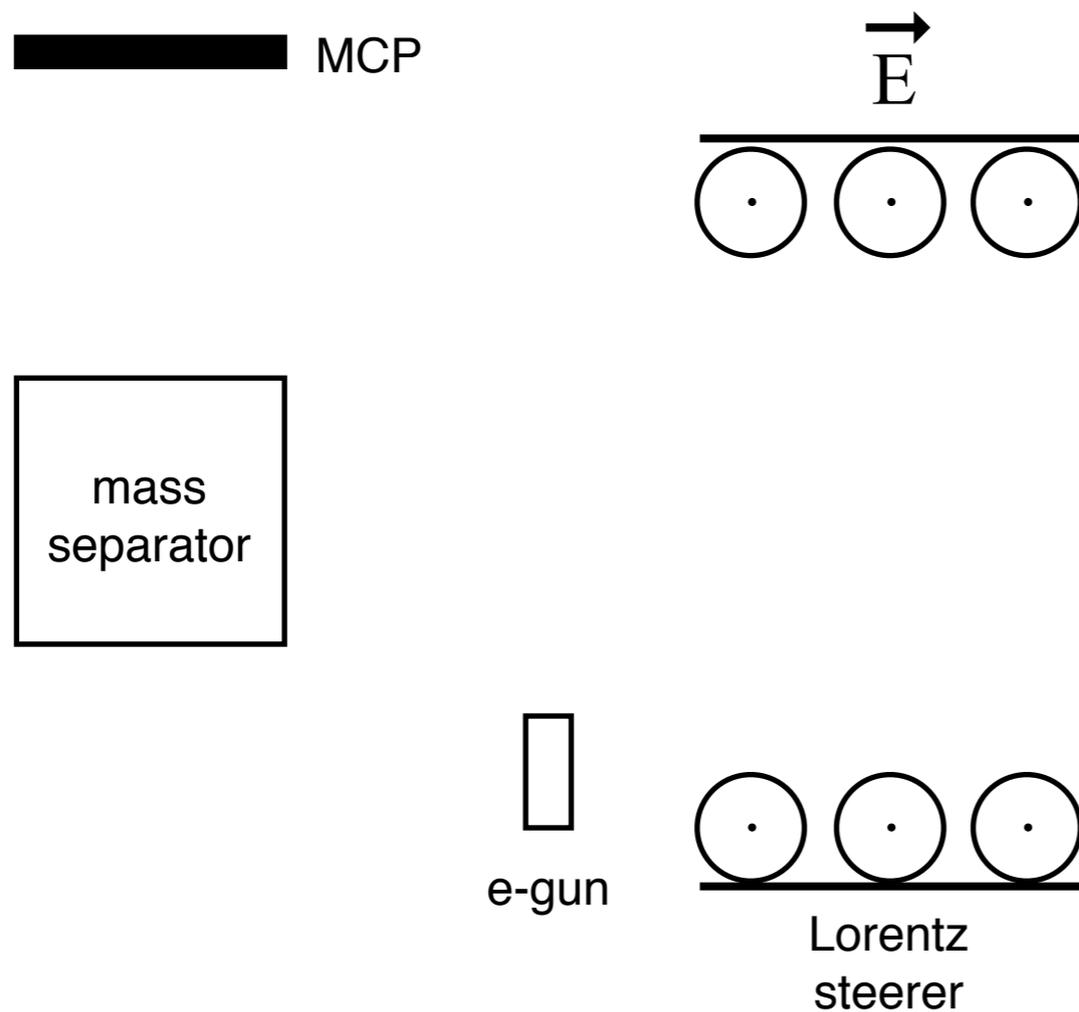
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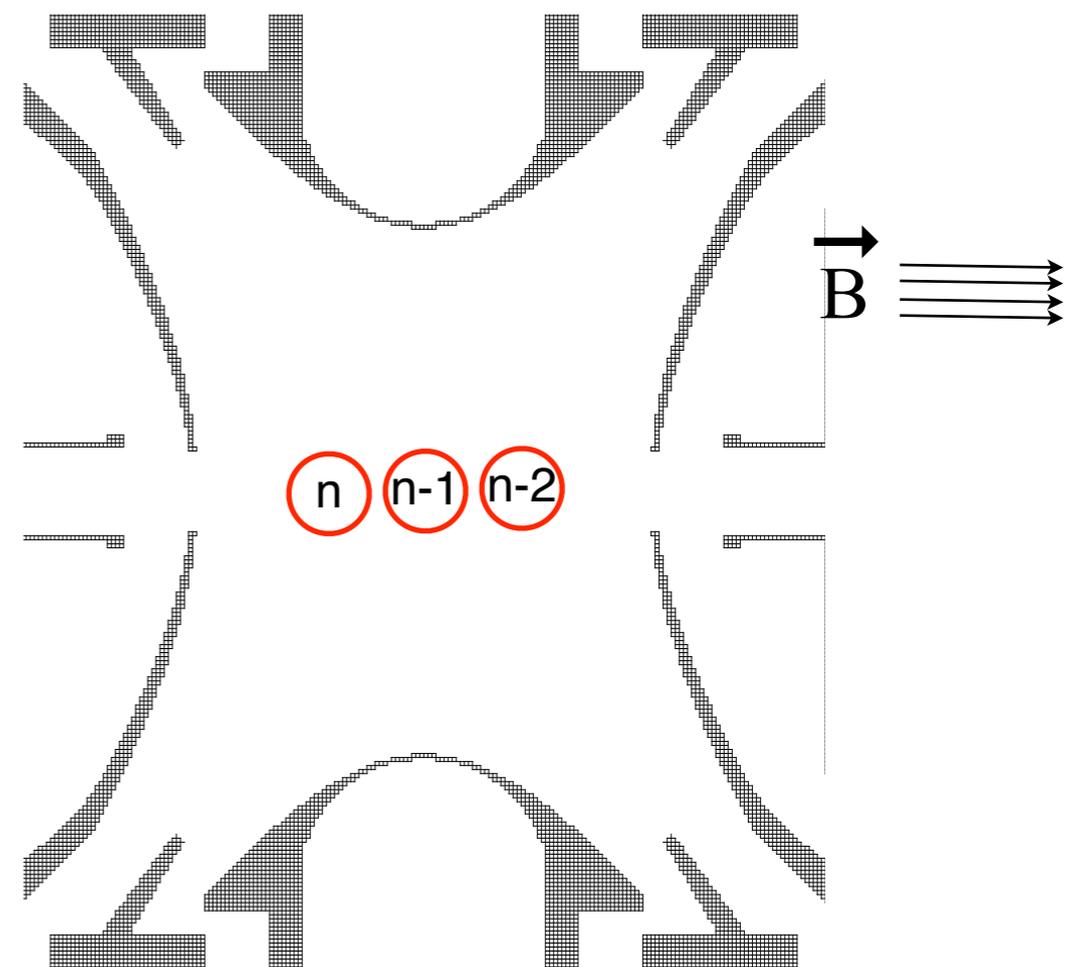
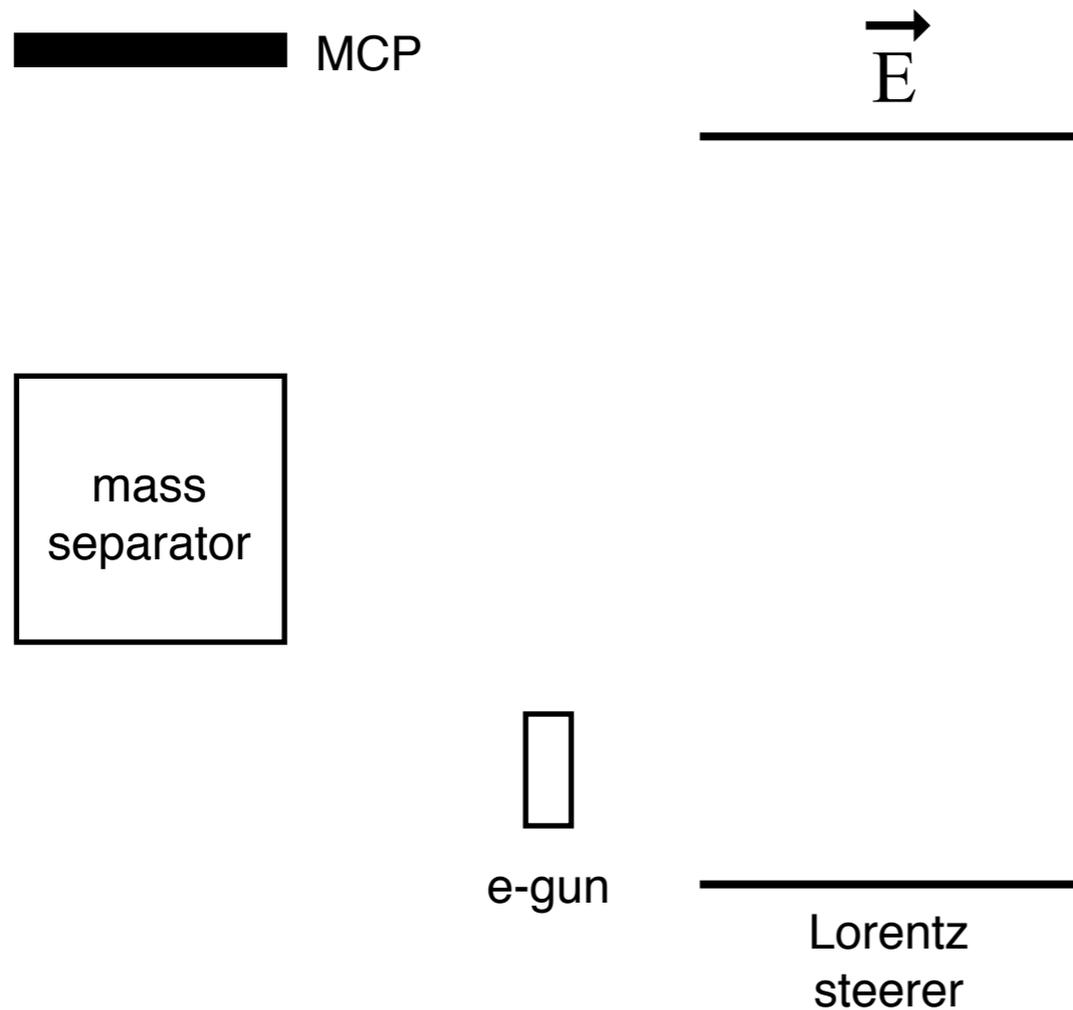
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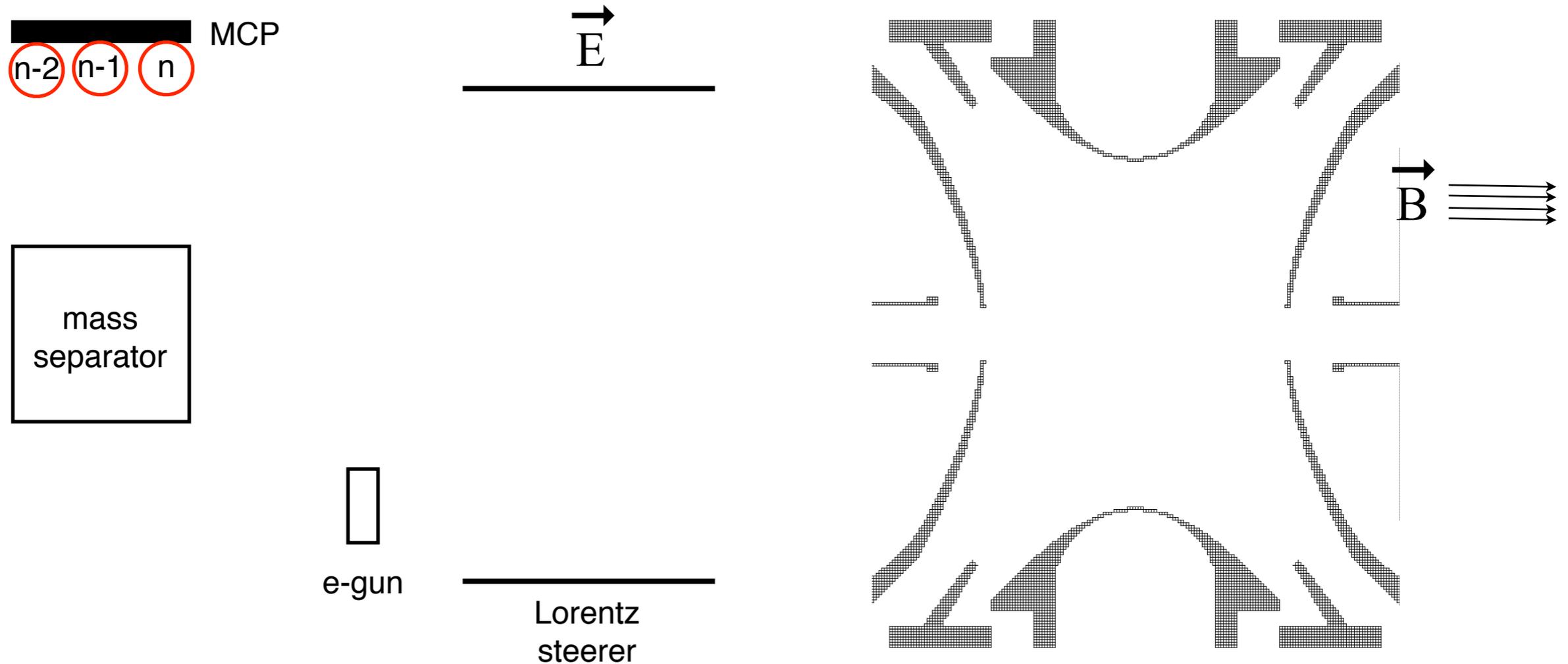
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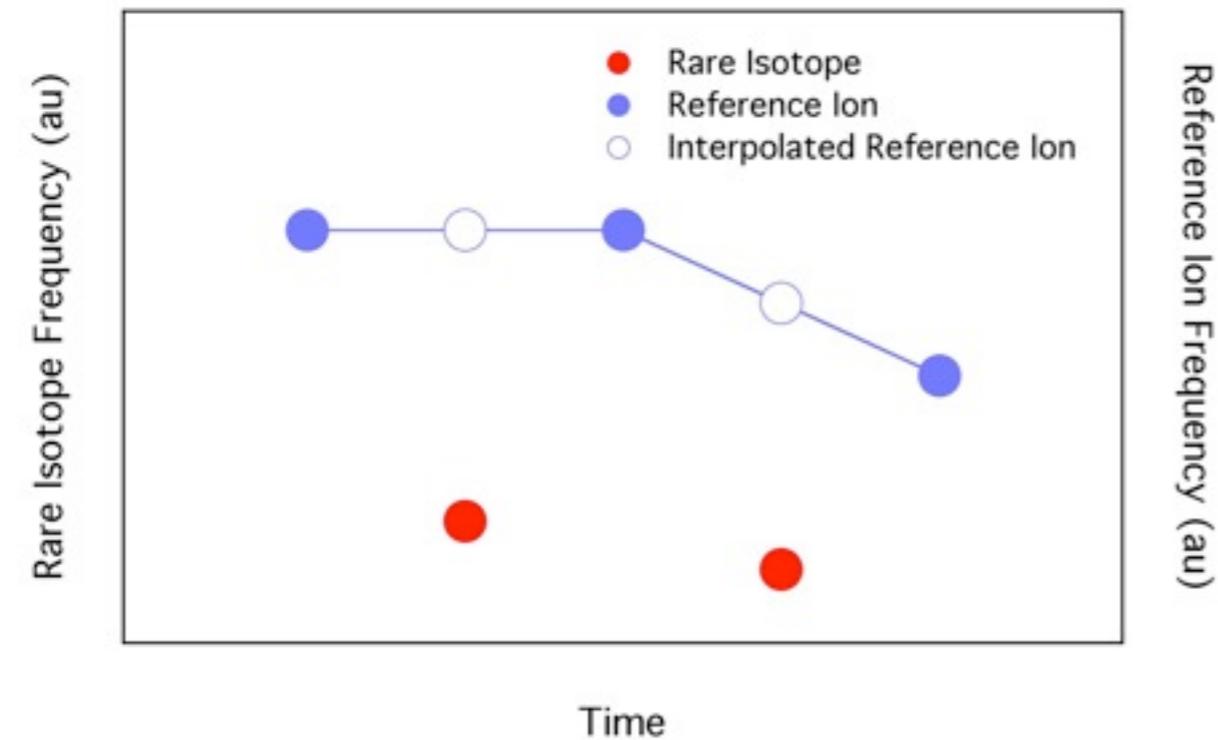
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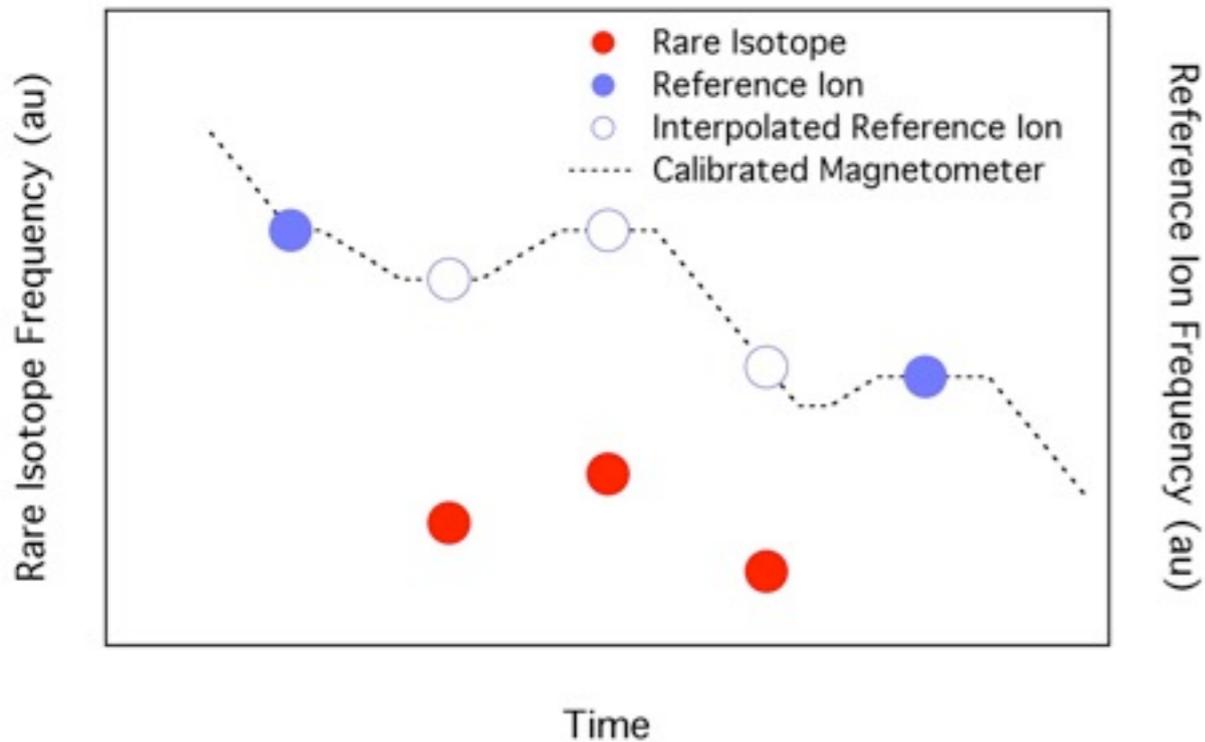
Mini Penning trap magnetometer



- Beam time is “wasted” on measuring reference ions
- Does not account for non-linear field drifts

Magnetic field is calibrated with a mass measurement of a reference ion before and after each RI ion measurement

Mini Penning trap magnetometer

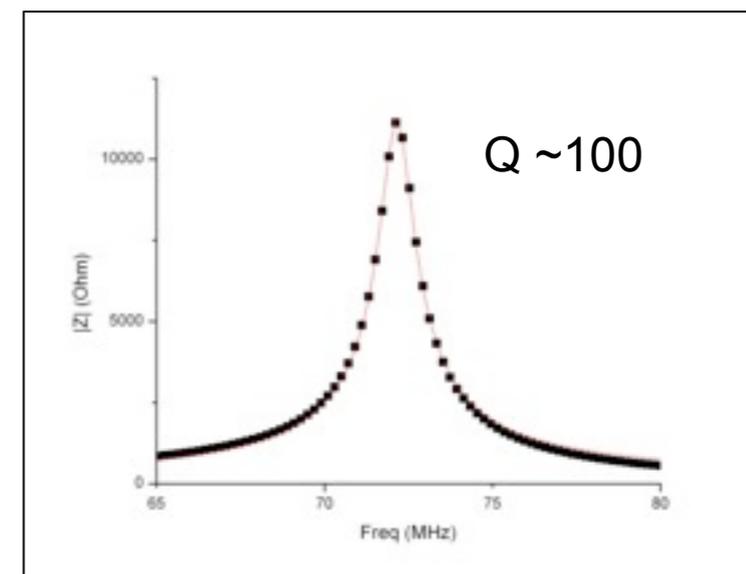
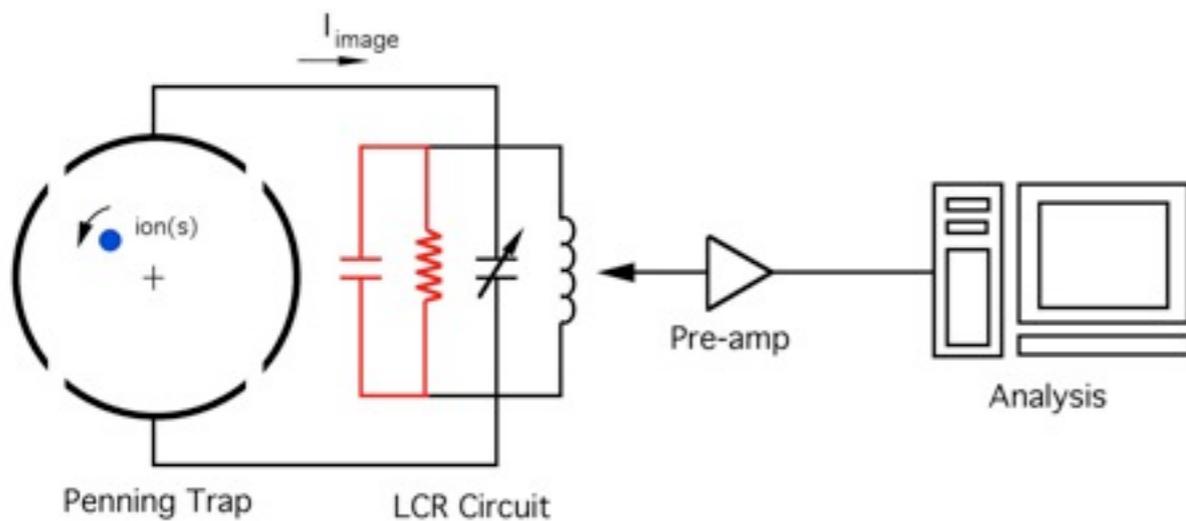


- Magnetic field is actively monitored during measurements
- Nonlinear drifts are systematically tracked

- Precision goal: $<10^{-8}$
- Aim to work with room temperature trap and detection

Trap size: ~ 0.5 cm
 Detector $Q \sim 100$
 # ions $\sim 100-1000$

$$\frac{S}{N} \sim Nq \left(\frac{\rho}{\rho_0} \right) \sqrt{\frac{\nu}{\Delta\nu}} \sqrt{\frac{Q}{kTC}}$$



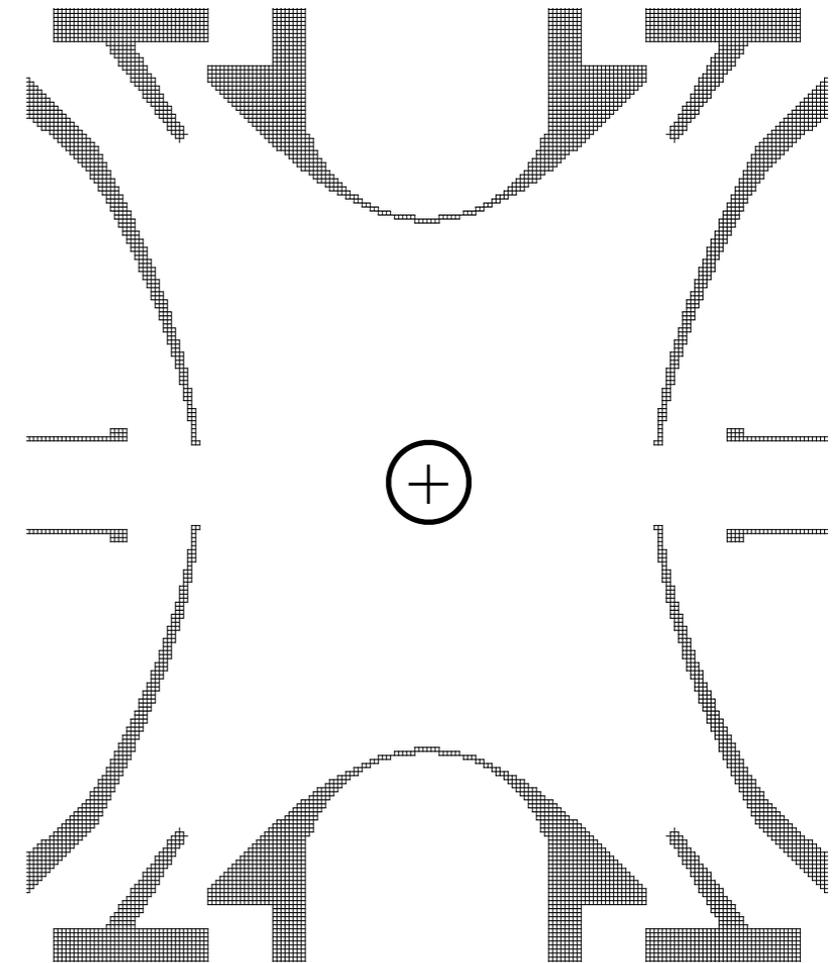
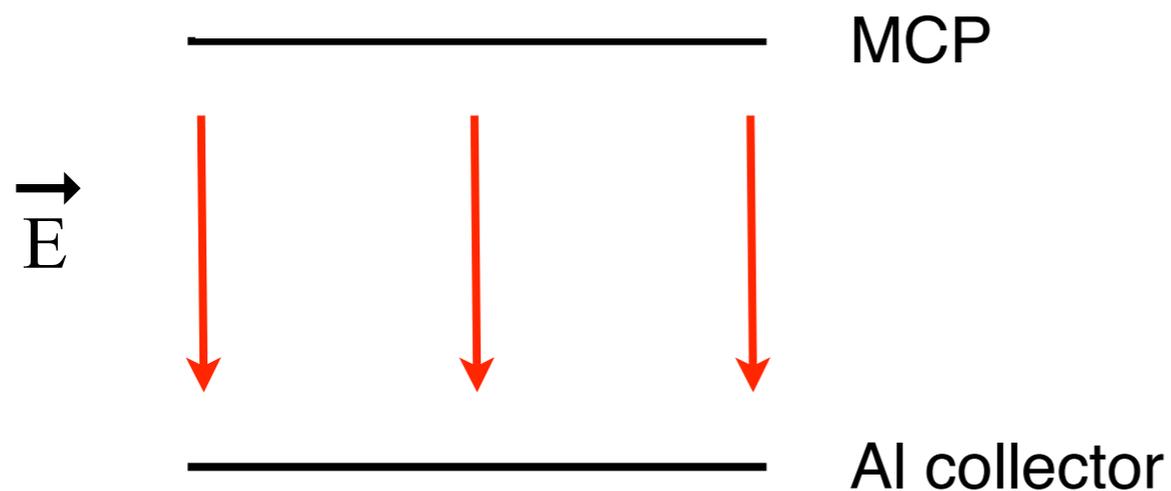
β -coincidence time of flight detection

background suppression for species with short half-life, low yield

$^{14}\text{Be} \sim 4 \text{ ms}$

$^{19}\text{C} \sim 46 \text{ ms}$

$^{70}\text{Kr} \sim 57 \text{ ms}$



N. R. Daly, *Review of Scientific Instruments* **31**, 264-7 (1960).

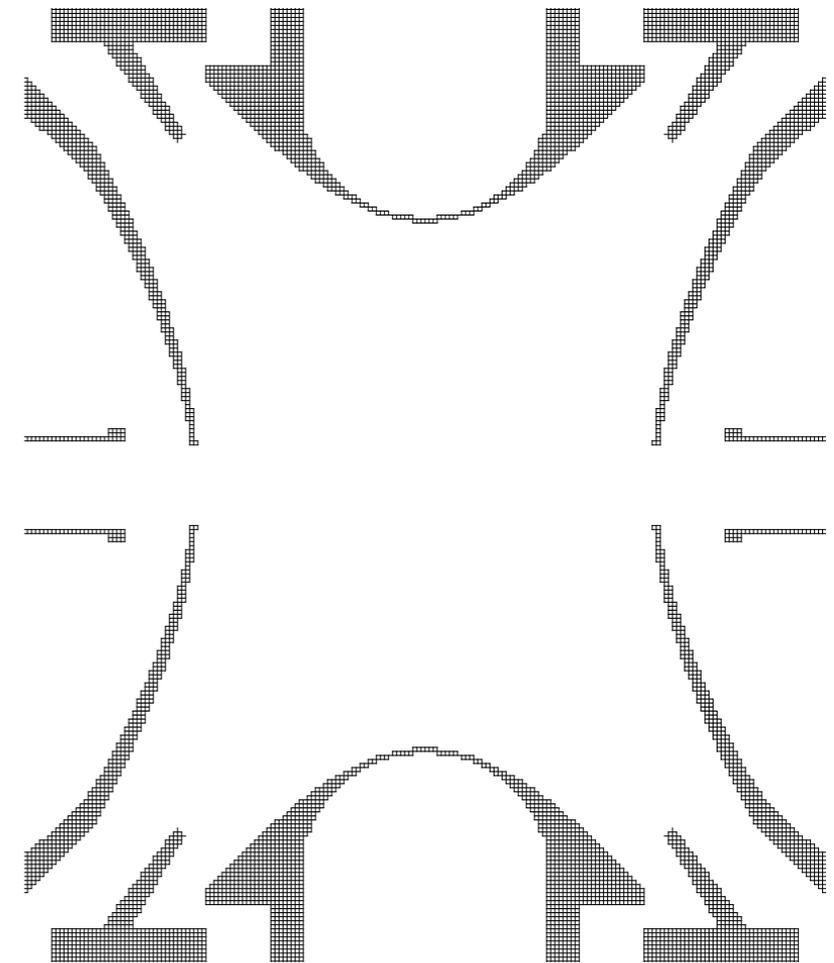
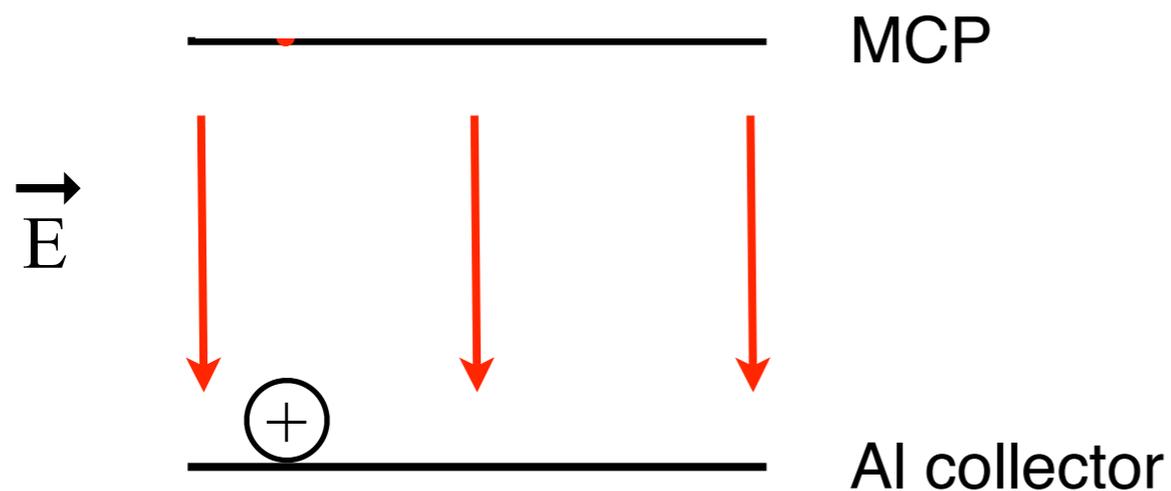
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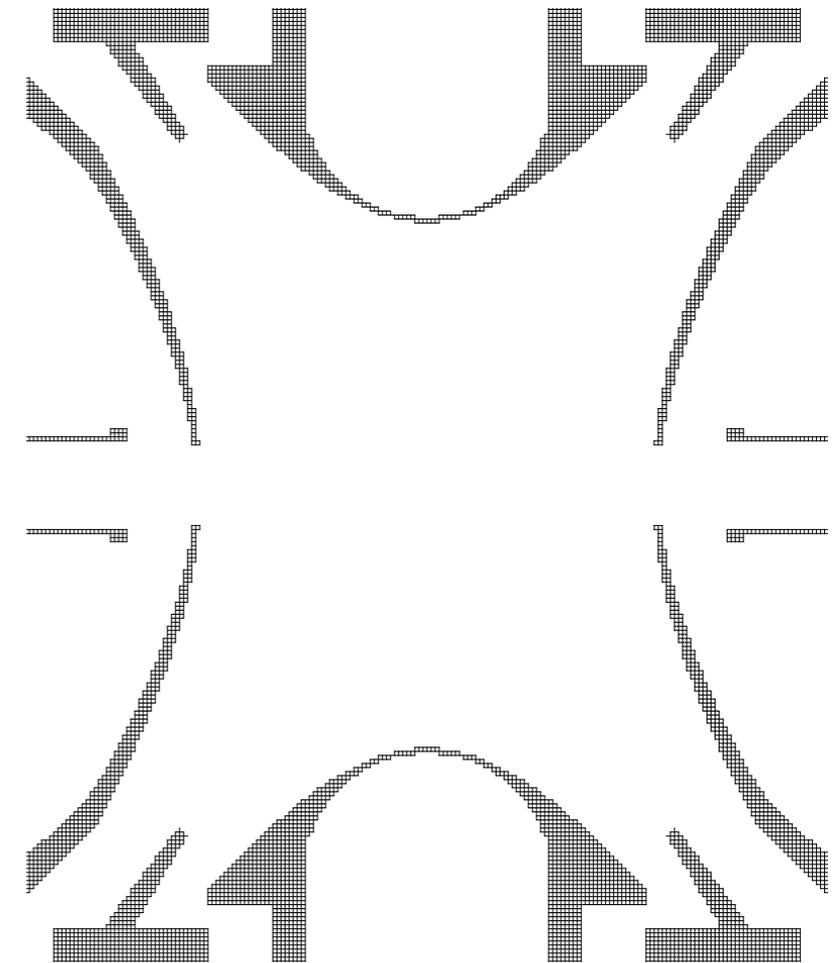
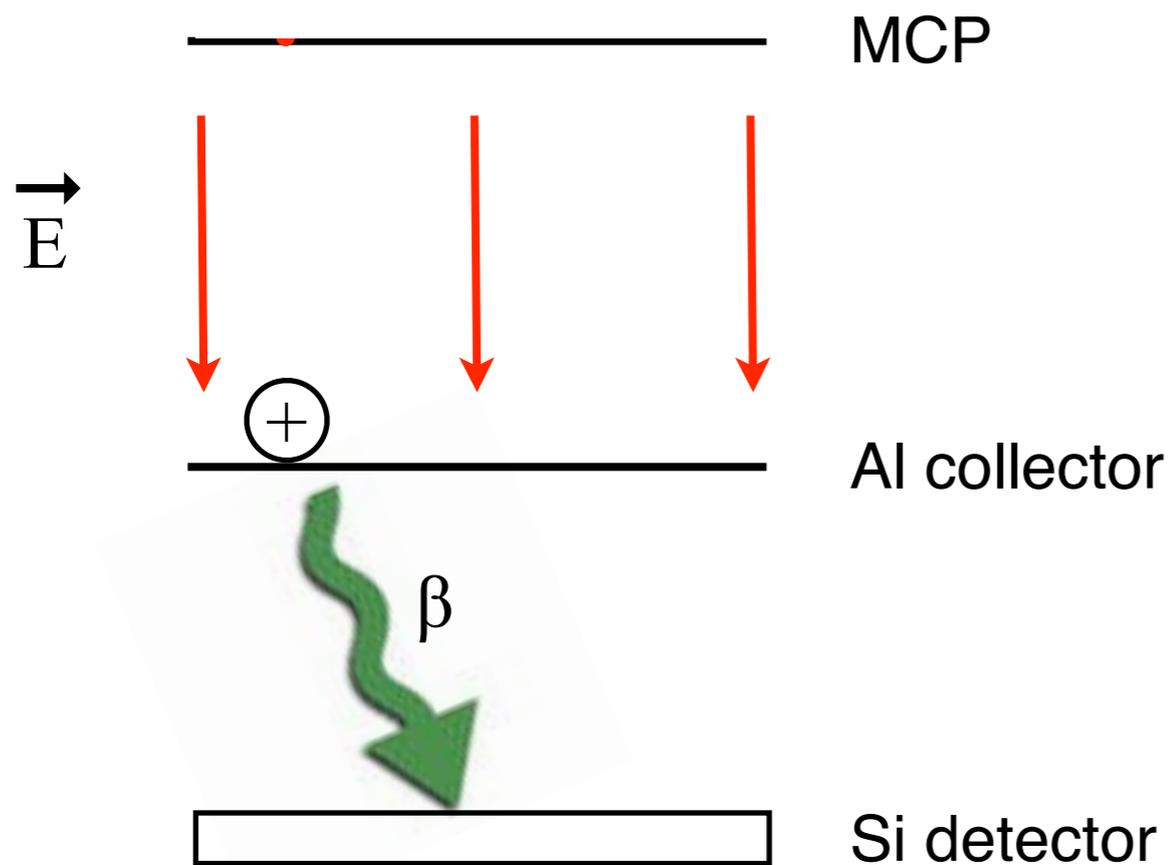
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3D cylindrical PIC code

Need to study space charge in traps for a variety of applications

Brute force: simulate ~ 1000 particles and scale Coulomb interaction for buffer gas cooling in Penning trap¹

3D cubic PIC: 10^6 ions in cubic domain used to study image-charge detection in FT-ICR²

Problems with these approaches

brute force: not very realistic, ignores image charges, gets expensive

cubic PIC: not a natural geometry for our traps, uniform cell volume

Benefits of 3D cylindrical PIC

- high cell density near origin
- natural boundary geometry
- ability to apply RF on boundaries
- add in physics (scattering, charge breeding, etc.)

Possible applications

- plasma evolution in trap
- side band or rotating wall cooling
- proton/electron cooling of HCl's
- thermalization in gas cells

1. D. Beck *et al.*, *Hyperfine Interactions*. **132**, 473-8 (2001).

2. E. N. Nikolaev *et al.*, *Rapid Communications in Mass Spectrometry* **21**, 3527-46 (2007).

